

P, C, and CP Transformation

For pseudoscalar mesons P and \overline{P} , the parity transformation implies

 $\mathbf{P}|P(\vec{p})\rangle = -|P(-\vec{p})\rangle, \qquad \mathbf{P}|\bar{P}(\vec{p})\rangle = -|\bar{P}(-\vec{p})\rangle. \qquad \mathbf{B}_{s}^{0} = \left(\overline{bs}\right) \quad \overline{B}_{s}^{0} = \left(b\overline{s}\right)$

Charge conjugation is a transformation that relates particles and antiparticles, leaving all space-time coordinates unchanged, i.e.

 $\mathbf{C} \left| P(\vec{p}) \right\rangle = \left| \bar{P}(\vec{p}) \right\rangle, \qquad \mathbf{C} \left| \bar{P}(\vec{p}) \right\rangle = \left| P(\vec{p}) \right\rangle.$

The combined transformation, **PC**, acts on the pseudoscalar meson states as follows:

 $B_d^0 = \left(\overline{b} d\right) \ \overline{B}_d^0 = \left(b\overline{d}\right)$

$$\mathbf{CP} \left| P(\vec{p}) \right\rangle = - \left| \bar{P}(-\vec{p}) \right\rangle, \qquad \mathbf{CP} \left| \bar{P}(\vec{p}) \right\rangle = - \left| P(-\vec{p}) \right\rangle.$$

For neutral mesons, P^0 and \overline{P}^0 , one can construct the *CP* eigenstates

$$|P_1^0\rangle = \frac{1}{\sqrt{2}} \left(|P^0\rangle - |\bar{P}^0\rangle \right), \qquad |P_2^0\rangle = \frac{1}{\sqrt{2}} \left(|P^0\rangle + |\bar{P}^0\rangle \right),$$

which obey

 $\mathbf{CP} \left| P_1^0 \right\rangle = \left| P_1^0 \right\rangle, \qquad \mathbf{CP} \left| P_2^0 \right\rangle = - \left| P_2^0 \right\rangle.$

Mixing

Effective Hamiltonian approximation: "dispersive" $i \frac{d}{dt} \begin{pmatrix} a(t) \\ b(t) \end{pmatrix} = H \begin{pmatrix} a(t) \\ b(t) \end{pmatrix}; P^0 = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \overline{P}^0 = \begin{pmatrix} 0 \\ 1 \end{pmatrix}; H_{ij} = M_{ij} - i\Gamma_{ij}/2$

From flavor to mass eigenstates $(P_L, P_H) \approx CP$ eigenstates (P_1, P_2) :

$$\left| P_{L}^{0} \right\rangle = p \left| P^{0} \right\rangle + q \left| \overline{P}^{0} \right\rangle = \frac{1}{\sqrt{1 + \left| \widetilde{\varepsilon} \right|^{2}}} \left(\widetilde{\varepsilon} \left| P_{1} \right\rangle + \left| P_{2} \right\rangle \right) \qquad \widetilde{\varepsilon} = \frac{p - q}{p + q}$$

$$\left| P_{H}^{0} \right\rangle = p \left| P^{0} \right\rangle - q \left| \overline{P}^{0} \right\rangle = \frac{1}{\sqrt{1 + \left| \widetilde{\varepsilon} \right|^{2}}} \left(\left| P_{1} \right\rangle + \widetilde{\varepsilon} \left| P_{2} \right\rangle \right) \qquad \left| q \right|^{2} + \left| p \right|^{2} = 1$$

Solving the eigenvalue equations and defining: $\Delta m = m_H - m_L$ $\Delta \Gamma = \Gamma_H - \Gamma_L$

$$\Delta m^2 - 1/4 \Delta \Gamma^2 = 4 |M_{12}|^2 - |\Gamma_{12}|^2$$
$$\Delta m \Delta \Gamma = 4 \Re e \left(M_{12} \Gamma_{12}^* \right)$$

q, p, Δm and $\Delta \Gamma$ for B_d and B_s



q, p, Δm and $\Delta \Gamma$ for B_d and B_s

$$M_{12} = -\frac{G_F^2 B_{B_d} f_{B_d}^2}{12\pi^2} m_B m_t^2 \eta_B V_{tb}^2 V_{td}^{*2} I\left(\frac{m_t^2}{m_W^2}\right), \quad I\left(\frac{m_t^2}{m_W^2}\right) = \begin{cases} 1., & m_t = 0\\ 0.5, & m_t = 175 GeV\\ 0.25, & m_t = \infty \end{cases}$$

$$\Gamma_{12} = \frac{G_F^2 B_{B_d} f_{B_d}^2}{8\pi} m_B^3 \left[-V_{tb} V_{td}^* + O\left(\frac{m_c^2}{m_b^2}\right) V_{cb} V_{cd}^* \right]^2$$

Where η_B with the account of NLO corrections ($\eta_B^{NLO}=0.55\pm0.01$) and $f_{B_d}\sqrt{B_{B_d}}=216\pm15 {
m MeV}$

In the SM M_{12} dominated by the top quark for B mesons: Γ_{12} few common on-shell states $\Gamma_{12}/M_{12} << 1$

$$\Rightarrow \Delta m \approx 2|M_{12}| \qquad \Delta \Gamma \approx \frac{2\Re e\left(M_{12}\Gamma_{12}^{*}\right)}{|M_{12}|} << \Delta m \qquad \frac{q}{p} = -\frac{\Delta m - i/2\Delta\Gamma}{2M_{12} - i\Gamma_{12}} \approx -\frac{|M_{12}|}{M_{12}}$$
$$CP-violating parameter: \qquad \delta = |p|^{2} - |q|^{2} = \langle P_{H} | P_{L} \rangle = \frac{2\Im m\left(M_{12}^{*}\Gamma_{12}\right)}{\left(\Delta m\right)^{2} + \left|\Gamma_{12}\right|^{2}} \approx 10^{-3}$$

Time evolution of neutral *B* mesons

Assuming CPT conservation

Time evolution of mass eigenstates:

$$\left| B_L^0(t) \right\rangle = e^{-t\Gamma_B/2} e^{-itM_B} e^{+it\Delta m_B/2} \left| B_L^0(0) \right\rangle$$
$$\left| B_H^0(t) \right\rangle = e^{-t\Gamma_B/2} e^{-itM_B} e^{-it\Delta m_B/2} \left| B_H^0(0) \right\rangle$$

Time evolution of initially (t=0) pure flavour eigenstates:

$$\begin{vmatrix} B_{phys}^{0}(t) \end{pmatrix} = h_{+}(t) \begin{vmatrix} B^{0} \end{pmatrix} + \frac{q}{p} h_{-}(t) \begin{vmatrix} \overline{B}^{0} \end{pmatrix}$$

$$h_{+}(t) = e^{-t\Gamma_{B}/2} e^{-itM_{B}} \cos(t \Delta m_{B}/2)$$

$$h_{-}(t) = i \left[e^{-t\Gamma_{B}/2} e^{-itM_{B}} \sin(t \Delta m_{B}/2) \right]$$

Time evolution of neutral **B** mesons

Flavour oscillations: for initially pure $B^0(t=0)$, probability for finding $B^0(\overline{B}^0)$ at time t, assuming |q/p|=1

$$|h_{\pm}(t)|^2 = \frac{1}{2} e^{-t\Gamma_B} \left[1 \pm \cos(t \,\Delta m_B)\right] \implies a_{mix}(t) = \cos(t \,\Delta m) = \cos(x\Gamma t)$$

Time-integrated ratio and time-integrated oscillation probability:

$$r = \frac{N(\overline{B}^{0})}{N(B^{0})} = \frac{\int_{0}^{\infty} dt \left|h_{-}(t)\right|^{2}}{\int_{0}^{\infty} dt \left|h_{+}(t)\right|^{2}} = \frac{x^{2}}{2+x^{2}}, \quad \chi = \frac{r}{1+r} = P(B^{0} \to \overline{B}^{0}), \quad x \equiv \frac{\Delta m}{\Gamma}$$

Observable by looking at self-flavour tagging semileptonic or hadronic decays! For example:

$$B^{0} \to D^{*-}l^{+}\nu \qquad B^{0} \to D^{*+}l^{-}\overline{\nu}$$
$$B^{0} \to D^{-}\pi^{+} \qquad \overline{B}^{0} \to D^{+}\pi^{-}$$
$$B^{0}_{s} \to D^{-}_{s}l^{+}\nu \qquad \overline{B}^{0}_{s} \to D^{+}_{s}l^{-}\overline{\nu}$$

Discovery $B\overline{B}$ oscillations

ARGUS Collaboration Observation of B – anti-B0 Mixing

Reconstructed Y(45) event

$$\begin{split} &\Upsilon(45) \to B^0 \bar{B^0} \to B_1^0 B_2^0 \\ &B_1^0 \to D_1^{*-} \mu_1^+ \nu_1, \ D_1^{*-} \to \bar{D^0} \pi_1^- \\ &B_2^0 \to D_2^{*-} \mu_2^+ \nu_2, \ D_1^{*-} \to D^- \pi^0 \end{split}$$

Time-integrated 21% mixing rate

- 25 (270) like (opposite) sign dilepton events
- 4.1 lepton-tagged semileptonic B decays

Integrated Y(45) Iuminosity 1983-87: • 103 pb⁻¹ ~ 110,000 B pairs



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$B_d \overline{B}_d$ Mixing at LHCb

$$A(t) = \frac{N^{\text{unmix}}(t) - N^{\text{mix}}(t)}{N^{\text{unmix}}(t) + N^{\text{mix}}(t)} = \cos(\Delta m_d t)$$



Mixing parameters



$B_s \overline{B}_s$ Mixing at LHCb

$$A(t) = \frac{N(B_s^0 \to D_s^- \pi^+, t) - N(\overline{B}_s^0 \to D_s^- \pi^+, t)}{N(B_s^0 \to D_s^- \pi^+, t) + N(\overline{B}_s^0 \to D_s^- \pi^+, t)},$$



CP Violation in B Decays



Classification of CP-violating effects
CPV in decay:

$$\begin{aligned} \left|\overline{A_{f}}/A_{f}\right| \neq 1 \\ A_{CP,f^{\pm}} &= \frac{\Gamma(P^{-} \rightarrow f^{-}) - \Gamma(P^{+} \rightarrow f^{+})}{\Gamma(P^{-} \rightarrow f^{-}) + \Gamma(P^{+} \rightarrow f^{+})} = \frac{\left|\overline{A_{f^{-}}}/A_{f^{+}}\right|^{2} - 1}{\left|\overline{A_{f^{-}}}/A_{f^{+}}\right|^{2} + 1} \\ CPV \text{ in mixing:} \\ A_{SL}(t) &= \frac{d\Gamma/dt(\overline{P}_{phys}^{0} \rightarrow l^{+}X) - d\Gamma/dt(P_{phys}^{0} \rightarrow l^{-}X)}{d\Gamma/dt(\overline{P}_{phys}^{0} \rightarrow l^{+}X) + d\Gamma/dt(P_{phys}^{0} \rightarrow l^{-}X)} = \\ &= \frac{1 - \left|q/p\right|^{4}}{1 + \left|q/p\right|^{4}} \end{aligned}$$

CPV in the interference decay-mixing:

 $\Im(\lambda_{f}) \neq 0$ For example: decays to CP eigenstates f_{CP} $\lambda_{f} \equiv \frac{q}{p} \frac{\overline{A}_{f}}{A_{f}}$ $A_{f_{CP}}(t) \equiv \frac{d\Gamma/dt \left(\overline{P}_{phys}^{0} \rightarrow f_{CP}\right) - d\Gamma/dt \left(P_{phys}^{0} \rightarrow f_{CP}\right)}{d\Gamma/dt \left(\overline{P}_{phys}^{0} \rightarrow f_{CP}\right) + d\Gamma/dt \left(P_{phys}^{0} \rightarrow f_{CP}\right)}$

Carter, Sanda PRL 45, 952 1980



Time-integrated "direct" CP asymmetry requires two amplitudes and $\delta \neq 0$:



Interference between mixing and decay to a CP eigenstate f_{CP} $\Rightarrow \Gamma(B^0_{phys}(t) \rightarrow f_{CP}) \neq \Gamma(\overline{B}^0_{phys}(t) \rightarrow f_{CP})$

Flavor-tagged time-dependent decay rates are different! they are governed by the "CP parameter":



Decay distributions $f_{+}(f)$ when tag = $B^{0}(\overline{B^{0}})$, pair-produced at Y(4S) $f_{CP,\pm}(\Delta t) = \frac{\Gamma}{4}e^{-\Gamma\Delta t}[1\pm S_{f_{CP}}\sin\Delta m_{d}\Delta t \mp C_{f_{CP}}\cos\Delta m_{d}\Delta t]$

Asymmetry

$$A_{f_{CP}}(\Delta t) = C_{f_{CP}} \cos(\Delta m_d \Delta t) - S_{f_{CP}} \sin(\Delta m_d \Delta t)$$

CP parameter

$$\lambda_{f_{CP}} = \eta_{f_{CP}} \frac{q}{p} \cdot \frac{\overline{A}_{\overline{f}_{CP}}}{A_{f_{CP}}}$$

$$\begin{split} C_{f_{CP}} &= \frac{1 - |\lambda_{f_{CP}}|^2}{1 + |\lambda_{f_{CP}}|^2} \\ S_{f_{CP}} &= \frac{-2 \ln \lambda_{f_{CP}}}{1 + |\lambda_{f_{CP}}|^2} \end{split}$$

For single decay amplitude = 0

 $=-\mathbf{Im}\lambda_{f_{CP}}$

Time Evolution of the Tagged $B^0(\overline{B}^0) \rightarrow B_{CP}$



For antisymmetric source of $B^0\overline{B}^0$, integrated CP asymmetry is zero: must do a time-dependent measurements

Golden Channel



KEKB asymmetric e⁺e⁻ collider



Time-Dependent CP Asymmetry Measurement



Detector Belle



Flavour tagging – dilution factor

$B^{\theta}B^{\theta} \rightarrow D^*l\nu$: reconstruction





CP asymmetry



$SIN(2\beta)$





2008 Nobel Prize in Physics

Makoto Kobayashi Toshihide Maskawa



for the discovery of the origin of the broken symmetry which predicts the existence of at least three families of quarks in nature

LHCb



LHCb

Time-dependent decay rate for
$$B_{phys}^{0} \rightarrow f$$
:

$$\frac{d\Gamma(B_{phys}^{0}(t) \rightarrow f)}{dt} = \left| \langle f | H | B_{phys}^{0}(t) \rangle \right|^{2} = \frac{d\Gamma(decay)}{dt}$$
"oscillation, then decay"
"oscillation, then decay"
"oscillation, then decay"

$$\left| (1 + \cos(\Delta m t)) | A_{f} |^{2} + (1 - \cos(\Delta m t)) | A_{f} |^{2} + (1 - \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_{f} |^{2} - (1 + \cos(\Delta m t)) | A_$$

Historical Remarks



Particle Identification System at Belle



Silicon Vertex Detector at Belle



Aerogel Cherenkov Counters



Electromagnetic Calorimeter



CsI(Tl) Crystals



 $\label{eq:light} \begin{array}{l} Light \ output - 5000 \ ph.el./MeV \\ Electronics \ noise \ \sigma{\sim}200 \ KeV \end{array}$



Detector Babar



- SVT: vertexing and tracking: crucial for Δt and low p_T tracks
- DCH: main tracking device, also dE/dx for particle ID
- DIRC: K- π separation > 3.4 σ for P < 3.5GeV/c
- EMC: very good energy resolution; electron ID, π^0 and γ reco.
- IFR: Muon and neutral hadrons (K⁰_L) ID

Silicon Vertex Tracker

- double-sided Si microstrip detectors
- 5 layers: 340 wafers, 150000 readout channels
- $20^{\circ} < \theta < 150^{\circ}$
- $\sigma_{\text{point}} \approx 10\text{-}15\,\mu\text{m} \text{ for the inner} \\ \text{layers}$





Silicon Vertex Tracker (Babar vs Belle)



- $\Delta z = z_{cp} z_{tag}$
 - $\Delta t \simeq \Delta z / (\gamma \beta c)$
- Interaction Point $\gg \Delta z$
- B flight-length in x-y: only $\sim 30\mu$
- C conservation in $\Upsilon(4S) \rightarrow B\overline{B}$
 - $\psi(t) = |B_1^0 > |\bar{B}_2^0 > -|\bar{B}_1^0 > |\bar{B}_2^0 >$

(one is B^0 and other is \bar{B}^0 at any time)

The other B provides time reference and flavor tagging at $\Delta t = 0$

| Parameters | BaBar | Belle |
|---|--|---|
| e ⁺ e ⁻ energy | 3.1 × 9 GeV | 3.5 × 8.5 GeV |
| γβ | 0.56 | 0.425 |
| Interaction point $(h \times v \times l)$ | $120\mu\text{m} \times 5\mu\text{m} \times 8.5\text{mm}$ | $80 \mu m \times 2 \mu m \times 3.4 mm$ |
| Typical Δz | 260µm | 200µ m |
| σ_z (CP-side) | $50 \mu m$ | 75µm |
| σ_z (tag-side) | $100 \sim 150 \mu\mathrm{m}$ | $140\mu\mathrm{m}$ |

DIRC



Identification Performance

Charged K identified by DIRC: Cerenkov angle DCH: dE/dx (p < 0.7 GeV/c)

Efficiency and purity measured on control samples (soft pion tag) $D^{*+} \rightarrow D^0\pi^+$, $D^0 \rightarrow K^-\pi^+$

> 3.4 $\sigma \pi/K$ separation up to \approx 3.5 GeV/c



"Golden Mode" Event



Luminosity



Observables: "direct" CP-violation

Time-integrated "direct" CP asymmetry ("CP violation in decay"):

$$A_{CP} = \frac{\Gamma(i \to f) - \Gamma(\bar{i} \to \bar{f})}{\Gamma(i \to f) + \Gamma(\bar{i} \to \bar{f})} = \frac{2|A_1||A_2|\sin\delta\sin\phi}{|A_1|^2 + |A_2|^2 + 2|A_1||A_2|\cos\delta\cos\phi}$$

- the only possibile CPV effect for charged mesons decays !
- requires at least two amplitudes and $\delta{\neq}0$

Reconstruction of B mesons



CPV Analysis: Time Distribution



 ω – mistag probability $R(\Delta t)$ - time-resolution function