

Generalized method of symmetric embeddings construction and its application to spacetimes of general relativity

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Outline

- Isometric embeddings
- Two-dimensional sphere example
- Bendable embeddings building
- $(2 + 1)$ -dimensional gravity examples
- Rotating BTZ example
- Conclusions

Isometric embeddings

Isometrically embedding a (pseudo)Riemannian manifold in a flat space means finding a surface in a flat space of a larger number of dimensions on which the induced metric will coincide with the metric of the original one.

Isometric embeddings

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However, even if an embedding in a flat space is found, it is not necessarily unique. In the case when the embeddings pass into each other with a continuous change of some parameter, then we are talking about the possibility of isometric bending of the embedding.

Two dimensional sphere example

$$ds^2 = d\theta^2 + \sin^2 \theta d\varphi^2. \quad (1)$$

2D sphere in \mathbb{R}^3 :

$$\begin{aligned} y^1 &= \sin \theta \cos \varphi, \\ y^2 &= \sin \theta \sin \varphi, \\ y^3 &= \cos \theta. \end{aligned} \quad (2)$$

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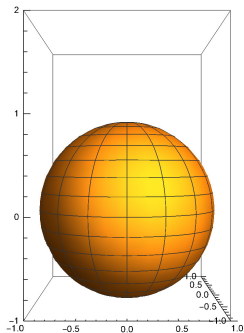
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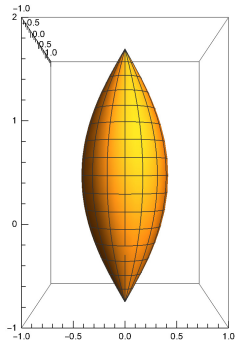
Bended sphere:

$$\begin{aligned} y^1 &= \frac{1}{a} \sin \theta \cos a\varphi, \\ y^2 &= \frac{1}{a} \sin \theta \sin a\varphi, \\ y^3 &= \int_0^\theta \sqrt{1 - \frac{\cos^2 \theta'}{a^2}} d\theta'. \end{aligned} \quad (3)$$

Two-dimensional sphere example



$$a = 1$$



$$a = 2$$

Building bendable embeddings

$$ds^2 = g_{ij}(x_1, \dots, x_n) dx^i dx^j + F^2(x_1, \dots, x_n) d\varphi^2 \quad (4)$$

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$$\begin{aligned} y^1 &= \sqrt{\xi} \frac{F}{a} \sin \sqrt{\xi} a \varphi, \\ y^2 &= \frac{F}{a} \cos \sqrt{\xi} a \varphi. \end{aligned} \quad (5)$$

$$\tilde{ds}^2 = g_{ij}(x_1, \dots, x_n) dx^i dx^j - \xi \frac{dF(x_1, \dots, x_n)^2}{a^2} \quad (6)$$

Building bendable embeddings

$$ds^2 = g_{ij}(x_1, \dots, x_n) dx^i dx^j + F^2(x_1, \dots, x_n) d\varphi^2 + G(x_1, \dots, x_n) dx^i d\varphi \quad (7)$$

Building bendable embeddings

$$ds^2 = g_{ij}(x_1, \dots, x_n) dx^i dx^j + F^2(x_1, \dots, x_n) d\varphi^2 + G(x_1, \dots, x_n) dx^i d\varphi \quad (7)$$

$$\begin{aligned} y^1 &= \sqrt{\xi} \sqrt{\frac{G(x_1, \dots, x_n)}{2bc}} \sin(\sqrt{\xi}(bx^i - c\varphi)), \\ y^2 &= \sqrt{\frac{G(x_1, \dots, x_n)}{2bc}} \cos(\sqrt{\xi}(bx^i - c\varphi)), \end{aligned} \quad (8)$$

Building bendable embeddings

$$ds^2 = g_{ij}(x_1, \dots, x_n) dx^i dx^j + F^2(x_1, \dots, x_n) d\varphi^2 + G(x_1, \dots, x_n) dx^i d\varphi \quad (7)$$

$$y^1 = \sqrt{\xi} \sqrt{\frac{G(x_1, \dots, x_n)}{2bc}} \sin(\sqrt{\xi}(bx^i - c\varphi)), \quad (8)$$
$$y^2 = \sqrt{\frac{G(x_1, \dots, x_n)}{2bc}} \cos(\sqrt{\xi}(bx^i - c\varphi)),$$

$$ds^2 = \tilde{g}_{ij}(x_1, \dots, x_n) dx^i dx^j + (F^2(x_1, \dots, x_n) + \eta G(x_1, \dots, x_n) \frac{c}{2b}) d\varphi^2 \quad (9)$$

(2 + 1)-dimensional gravity examples

$\Lambda > 0$ ¹:

$$ds^2 = \cos^2 \theta dt^2 - \frac{1}{\Lambda} (d\theta^2 + \sin^2 \theta d\varphi^2), \quad (10)$$

где $t \in (-\infty, \infty)$, $\theta \in [0, \pi)$, $\varphi \in [0, 2\pi(1 - 4GM))$.

¹S. Desert and R. Jekeyw. Three-dimensional cosmological gravity: dynamics of constant curvature. *Annals of Physics*, 1984

(2 + 1)-dimensional gravity examples

$$\begin{aligned}y^1 &= \frac{1}{a\sqrt{\Lambda}} \cos \theta \sinh a\sqrt{\Lambda}t, \\y^2 &= \frac{1}{a\sqrt{\Lambda}} \cos \theta \cosh a\sqrt{\Lambda}t, \\y^3 &= \frac{1}{b\sqrt{\Lambda}} \sin \theta \sin b\varphi, \\y^4 &= \frac{1}{b\sqrt{\Lambda}} \sin \theta \cos b\varphi. \\y^5 &= \frac{1}{\sqrt{\Lambda}} \int_0^\theta \sqrt{1 - \frac{\sin^2(\theta')}{a^2} - \frac{\cos^2(\theta')}{b^2}} d\theta'\end{aligned}\tag{11}$$

, where $\varphi \in [0, \frac{2\pi}{b})$, $b = 1/(1 - 4GM) \geq 1$, $a \geq 1$.

(2 + 1)-dimensional gravity examples

Let $a = b$

$$\Rightarrow y^5 = \frac{1}{\sqrt{\Lambda}} \sqrt{1 - \frac{1}{b^2} \theta} \quad (12)$$

(2 + 1)-dimensional gravity examples

Magnetic monopole in (2+1) dimensions², $\Lambda = -1/l^2$:

$$ds^2 = \frac{1}{l^2}(x^2 + \bar{r}^2 - r_+^2)dt^2 - \frac{l^2 x^2 dx^2}{(x^2 + \bar{r}^2 - r_+^2) \left(x^2 + Q_m^2 \ln \left[1 + \frac{x^2}{\bar{r}^2 - r_+^2} \right] \right) - \left(x^2 + Q_m^2 \ln \left[1 + \frac{x^2}{\bar{r}^2 - r_+^2} \right] \right) d\varphi^2} \quad (13)$$

where Q_m, M, \bar{r} – some parameters, $r_+^2 = Ml^2$

$t \in (0, \infty), x \in (0, \infty), \varphi \in (0, 2\pi T_\varphi]$ and $T_\varphi = \frac{\exp(\bar{r}^2/2Q_m^2)}{1 + Q_m^2 \exp(\bar{r}^2/2Q_m^2)}$

²Hirschmann E. W., Welch D. L. Magnetic solutions to 2+ 1 gravity //Physical Review D. – 1996. – T. 53. – №. 10. – C. 5579.

$$\begin{aligned}
y^1 &= \frac{1}{al} \sqrt{x^2 + \bar{r}^2 - r_+^2} \sinh at, \\
y^2 &= \frac{1}{al} \sqrt{x^2 + \bar{r}^2 - r_+^2} \cosh at, \\
y^3 &= \frac{1}{b} \sqrt{x^2 + Q_m^2 \ln \left[1 + \frac{x^2}{\bar{r}^2 - r_+^2} \right]} \sin b\varphi \\
y^4 &= \frac{1}{b} \sqrt{x^2 + Q_m^2 \ln \left[1 + \frac{x^2}{\bar{r}^2 - r_+^2} \right]} \cos b\varphi \\
y^5 &= \int lx(x^2 + \bar{r}^2 - r_+^2)^{-1/2} (x^2 + Q_m^2 \ln[1 + \frac{x^2}{\bar{r}^2 - r_+^2}])^{-1/2} dx, \\
y^6 &= \int \frac{1}{abl} \frac{x}{f(x)} \sqrt{\frac{G(x)}{x^2 + Q_m^2 \ln[1 + \frac{x^2}{\bar{r}^2 - r_+^2}]} } dx, \\
G(x) &= a^2 l^2 f(x)^2 + (b^2 x^2 + 2Q_m^2 a^2 l^2 + Q_m^2 b^2 \ln[1 + \frac{x^2}{\bar{r}^2 - r_+^2}]) f(x) + Q_m^2 l^2 a^2, \\
f(x) &= x^2 + \bar{r}^2 - r_+^2
\end{aligned} \tag{14}$$

Rotating BTZ example

Rotating BTZ in Eddington-Finkelstein coordinates³:

$$ds^2 = \left(-M + \frac{r^2}{l^2} \right) dv^2 - 2dvdr + Jdv d\bar{\varphi} - r^2 d\bar{\varphi}^2 \quad (15)$$

³Sheykin A. A., Markov M. V., Paston S. A. Global embedding of BTZ spacetime using generalized method of symmetric embeddings construction //Journal of Mathematical Physics. – 2021. – T. 62. – №. 10.

Rotating BTZ example

$$\begin{aligned}y^1 &= \sqrt{\frac{1}{a}} \sin\left(\frac{aJ}{2}v + \bar{\varphi} - \frac{2}{J}r\right), \\y^2 &= \sqrt{\frac{1}{a}} \cos\left(\frac{aJ}{2}v + \bar{\varphi} - \frac{2}{J}r\right), \\y^3 &= \sqrt{r^2 + \frac{1}{a}} \cos\left(\bar{\varphi} - \frac{2}{J} \frac{\arctan(\sqrt{ar})}{\sqrt{a}}\right), \\y^4 &= \sqrt{r^2 + \frac{1}{a}} \sin\left(\bar{\varphi} - \frac{2}{J} \frac{\arctan(\sqrt{ar})}{\sqrt{a}}\right), \\y^5 &= bv, \\y^6 &= \frac{1}{c} \sqrt{b^2 - M + \frac{r^2}{l^2} - \frac{aJ^2}{4}} \sin(cv), \\y^7 &= \frac{1}{c} \sqrt{b^2 - M + \frac{r^2}{l^2} - \frac{aJ^2}{4}} \cos(cv),\end{aligned}\tag{16}$$

Conclusions

- Has been developed a method that makes it possible to construct isometrically bendable embeddings of sufficiently symmetric spaces.
- It turns out that the resulting method is also useful for the problem of constructing explicit embeddings in general.
- The main result from the point of view of the problem of constructing explicit embeddings is the global embedding of a BTZ black hole into a 7-dimensional flat space.